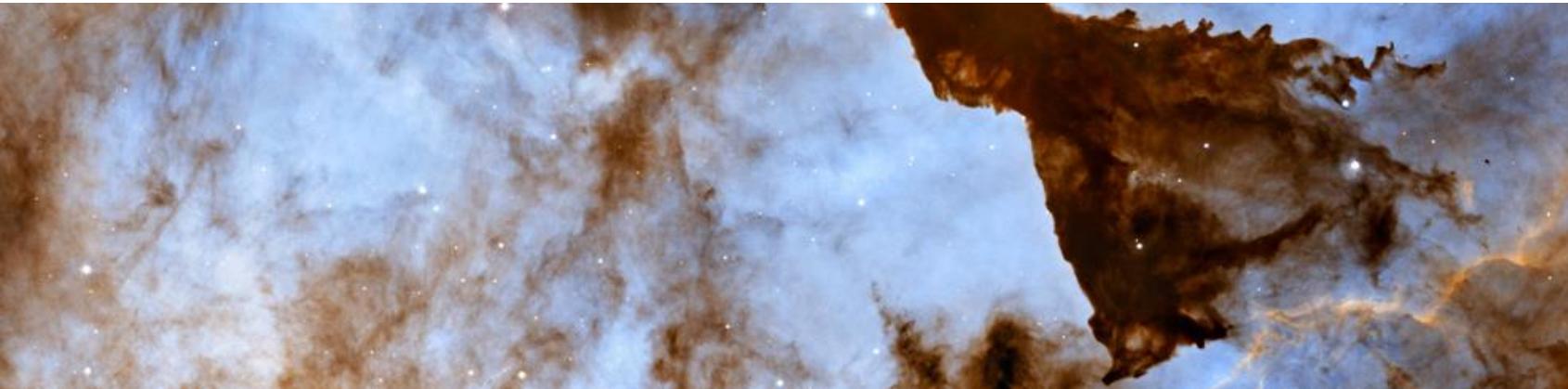


HLADNI MOLEKULSKI OBLACI



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Osobine

- najgušća vrsta MZM
- objekti sa sopstvenom gravitacijom
- veoma turbulentni
- prisutno magnetno polje
- akrecija na česticama prašine

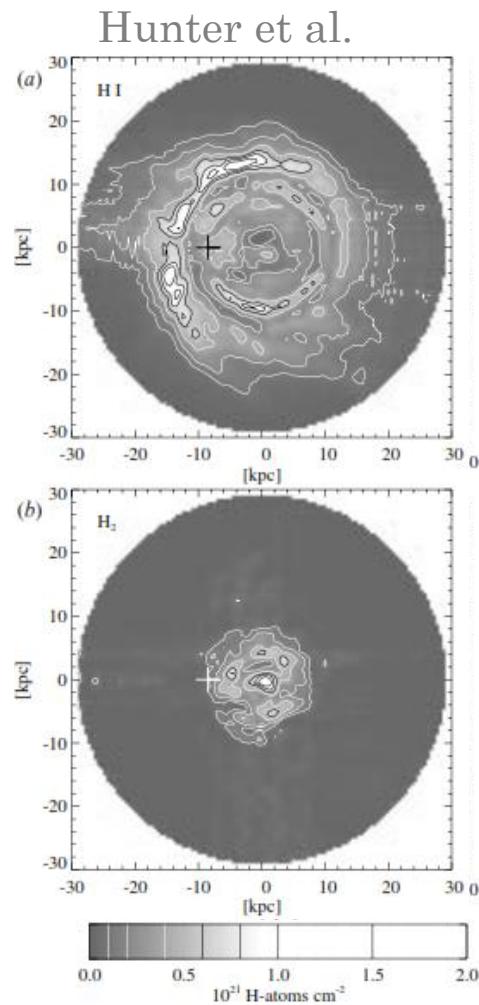


dominantan
fizički proces?

	<i>Veličina</i>	<i>Masa i gustina</i>	<i>Predstavnici</i>	<i>Karakteristike</i>
Džinovski molekulski oblaci	Oko 50–120 pc	$\approx 10^6 M_{\odot}$; $n \approx 10^2$ – 10^3 čestica po cm^{-3}	Orionov džinovski molekulski oblak, Maglina Orao	Detektovano je preko 100 različitih molekula
Mali oblaci, globule	Oko 1 sg	2 – $100 M_{\odot}$	Bokove globule	Okolina u H II regionima; sadrže H_2 , CO i He i silikatne čestice
Difuzni vlaknasti oblaci	Do 100 pc	30 čestica po cm^{-3}	Infracrveni cirusi	Detektovani na velikim galaktičkim širinama - gusta, amonijačna jezgra.

Prostorna raspodela MC

- zauzima manje od 1% zapremine međuzvezdanog prostora
 - između 3.5 i 7.5 kpc od centra Galaksije
 - disk debljine 50-75 pc
- položaj usko povezan sa spiralnim granama
 - protivi se teoriji disocijacije ($t < 10^6$ god)



CO mapa naše Galaksije

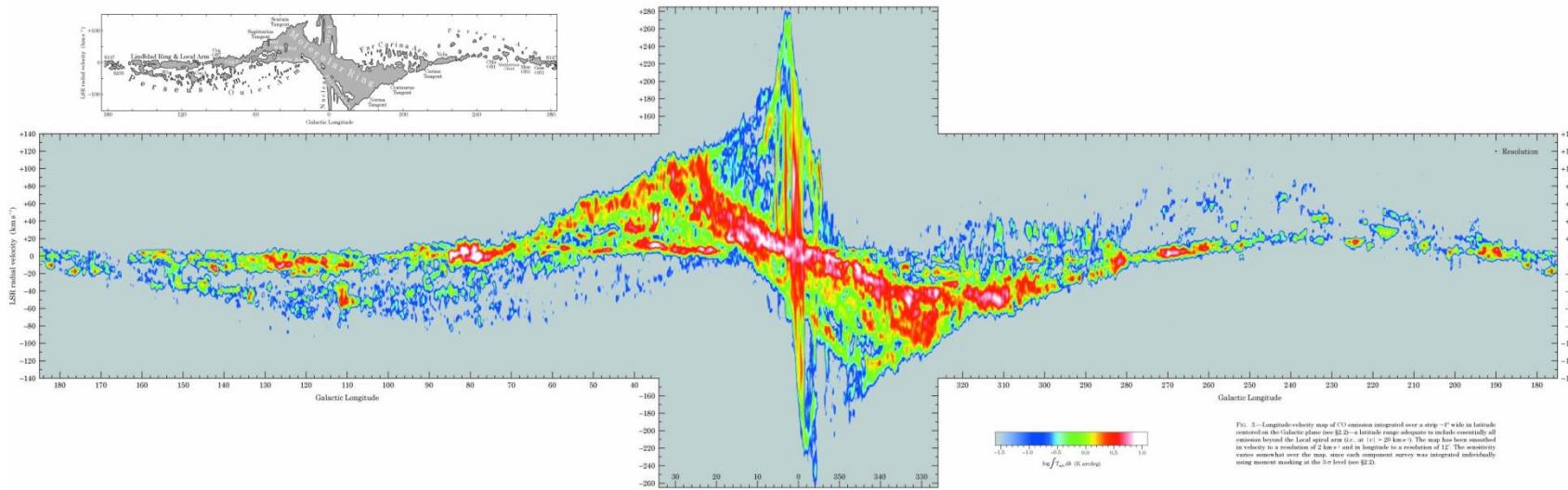
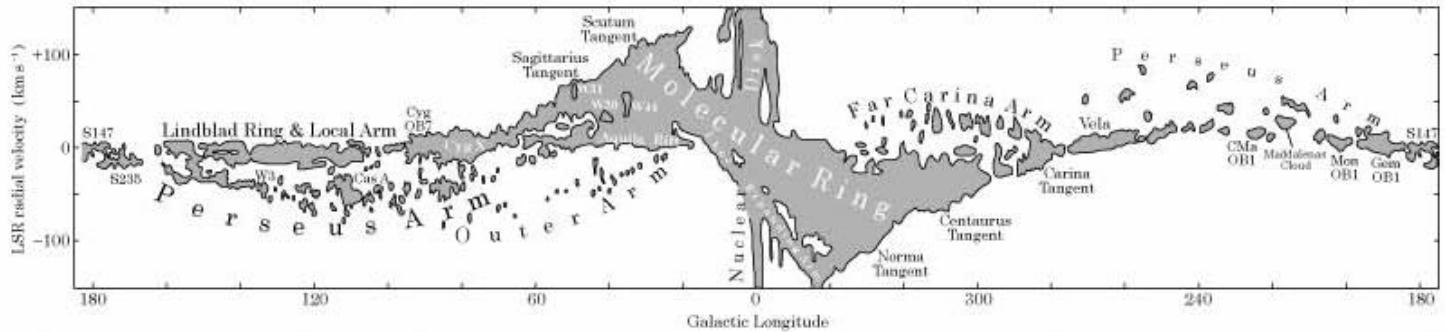
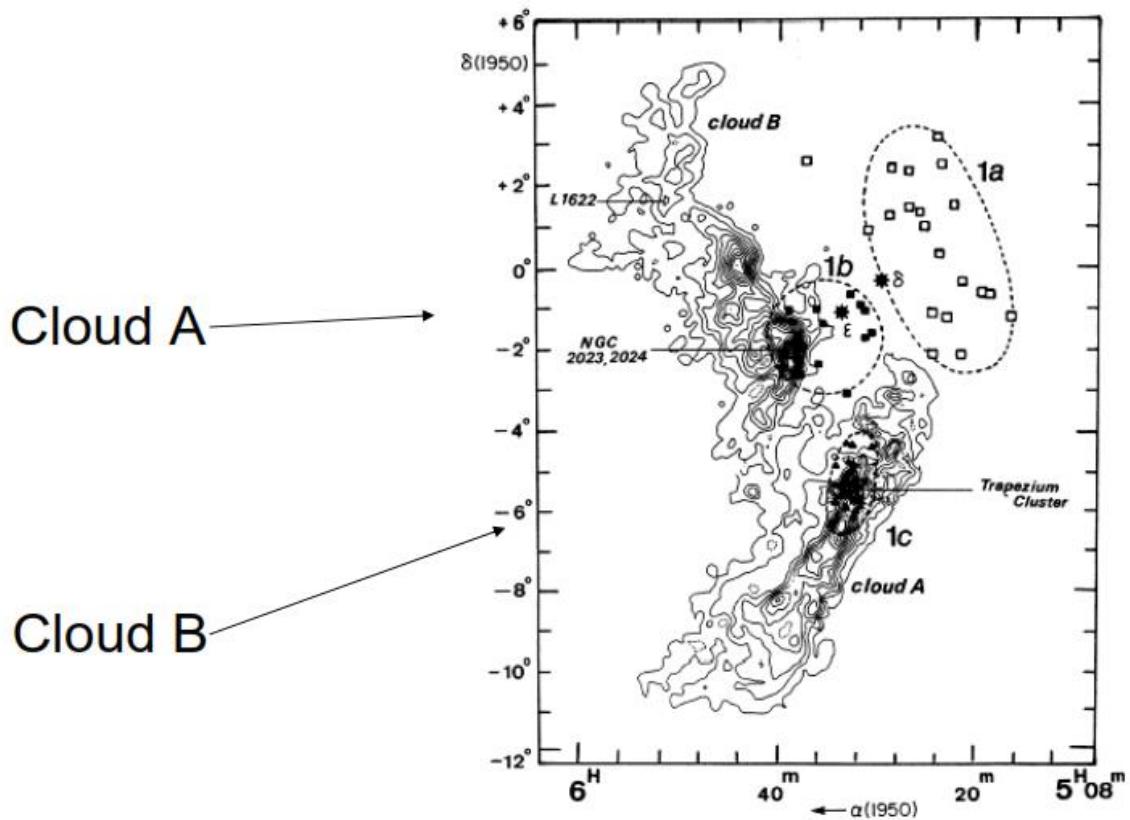
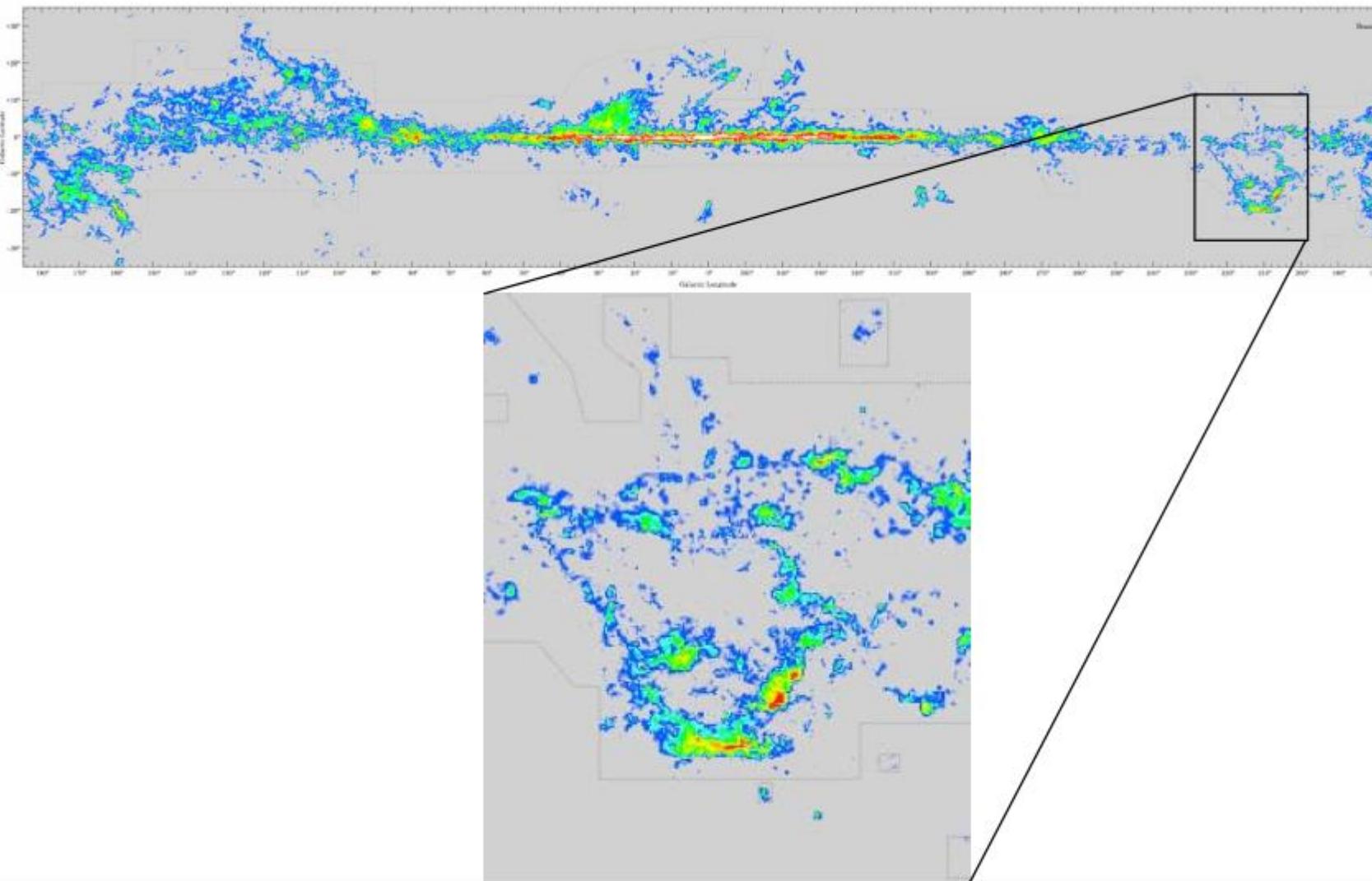


FIG. 3.—Longitude-velocity map of CO emission integrated over a strip $\sim 4^\circ$ wide in Latitude centered on the Galactic plane (see §2.2)—a latitude range adequate to include essentially all emission beyond the Local spiral arm (i.e., at $|v| > 20 \text{ km s}^{-1}$). The map has been smoothed in velocity to a resolution of $\pm 2 \text{ km s}^{-1}$ and in longitude to a resolution of $12'$. The sensitivity varies somewhat over the map, since each component survey was integrated individually using moment masking at the 3σ level (see §2.2).

Primer oblaka Orion A i B

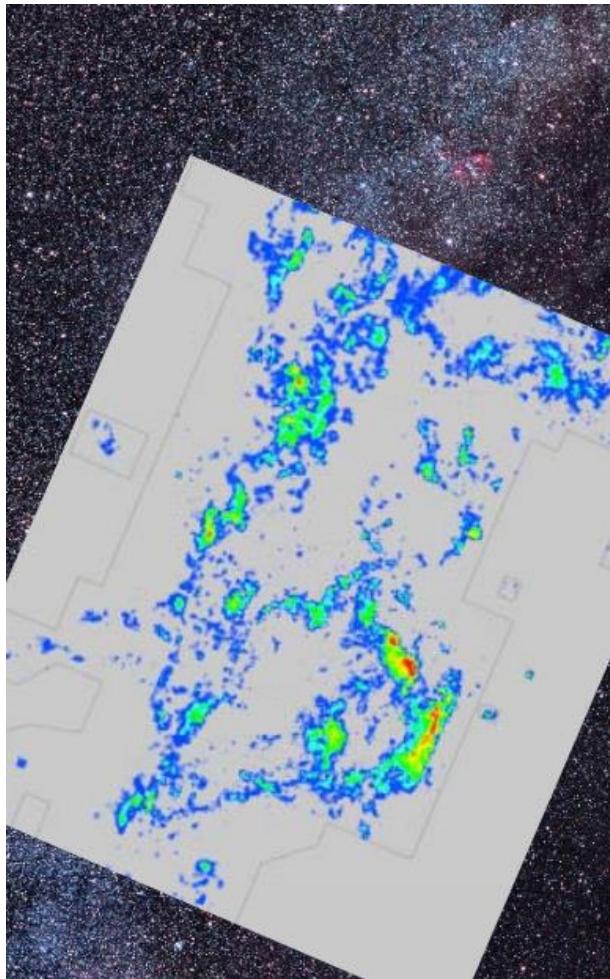


Blitz & Williams 1999

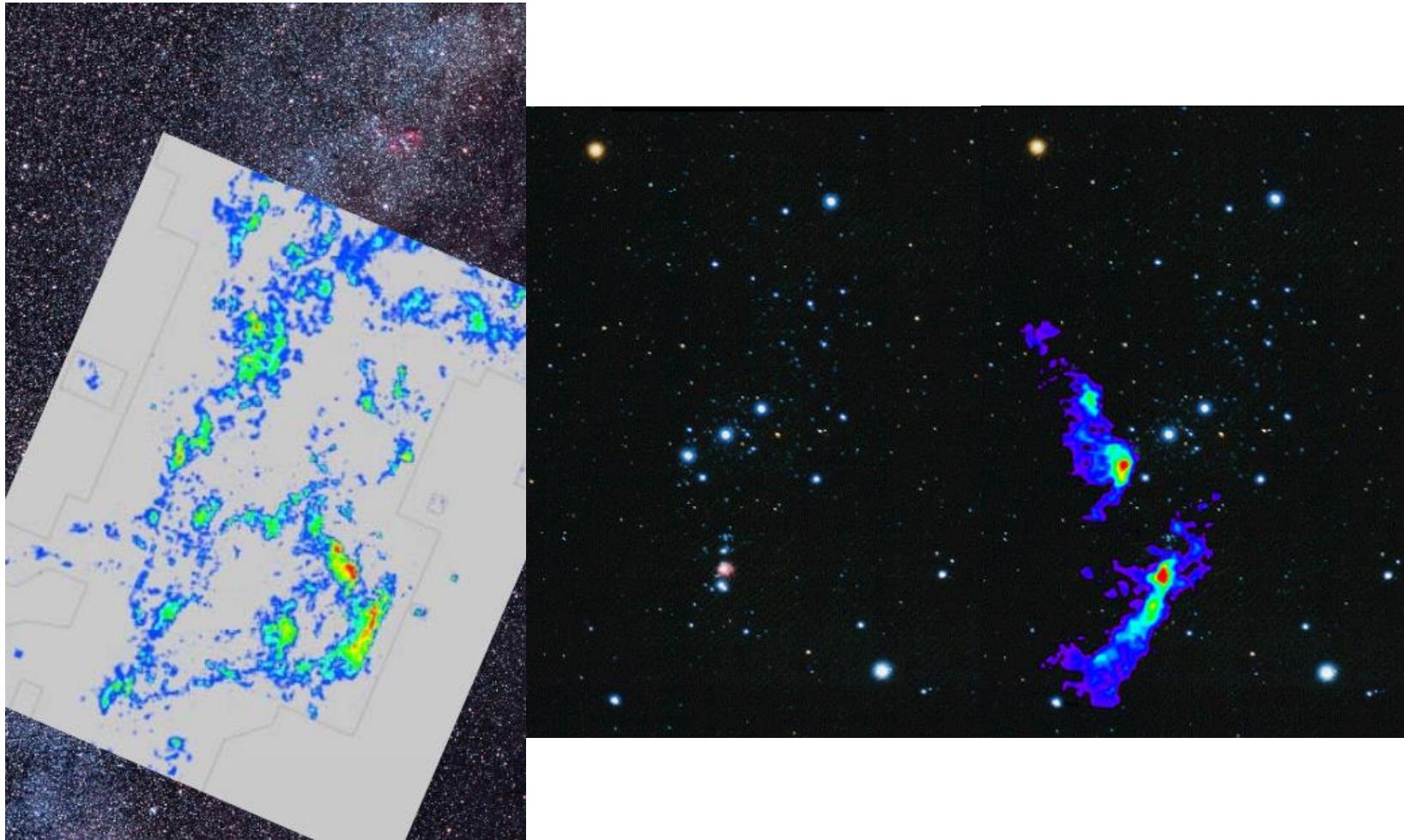


Dame et al. 2001

Simnci u vidljivom i mape CO



Simnici u vidljivom i mape CO



CO mape drugih galaksija

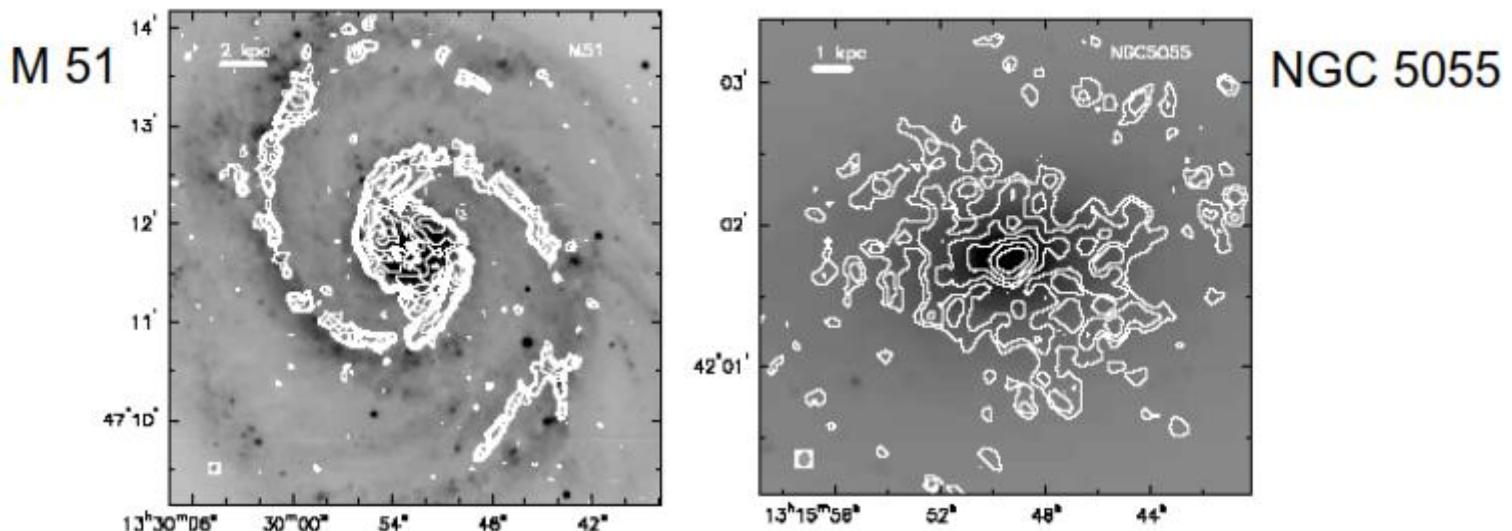
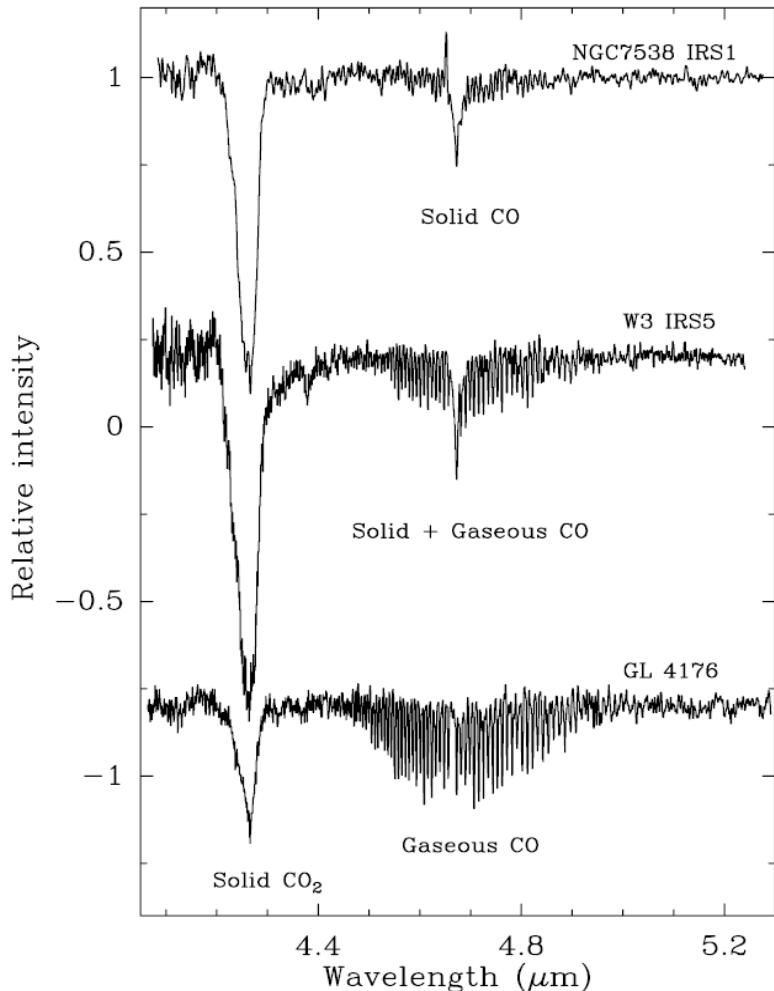


Figure 1. Left: CO in M51 from the BIMA Survey of Nearby Galaxies (SONG), overlaid on an optical image of the galaxy. The image contains zero spacing data and thus samples all of the spatial frequencies to the resolution limit (shown by the small box in the lower left). The CO is very strongly concentrated to the spiral arms and lies upstream of most of the ionized gas. Right: BIMA SONG image of the CO in NGC 5055 overlaid on an optical image showing that the galaxy is devoid of large-scale spiral structure at visible wavelengths. Some of the off-nuclear CO peaks are associated with weak spiral arms seen in the near infrared. The difference in CO morphology between these two galaxies is rather striking.

Primer spektra CO

- Apsorpcija na spektrima mladih zvezda
 - posmatranja ka tri zvezde
 - spektar različitih agregatnih stanja



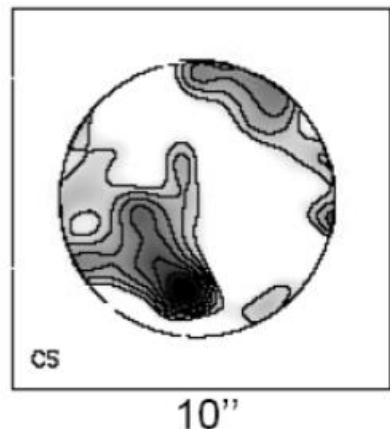
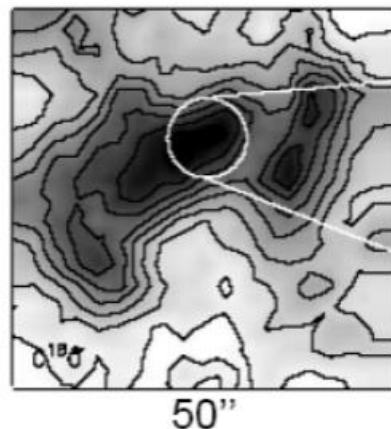
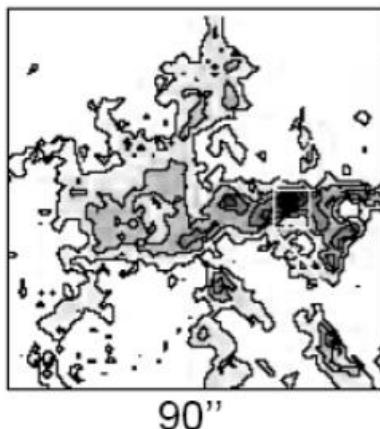
Spisak molekula detektovanih do 2003

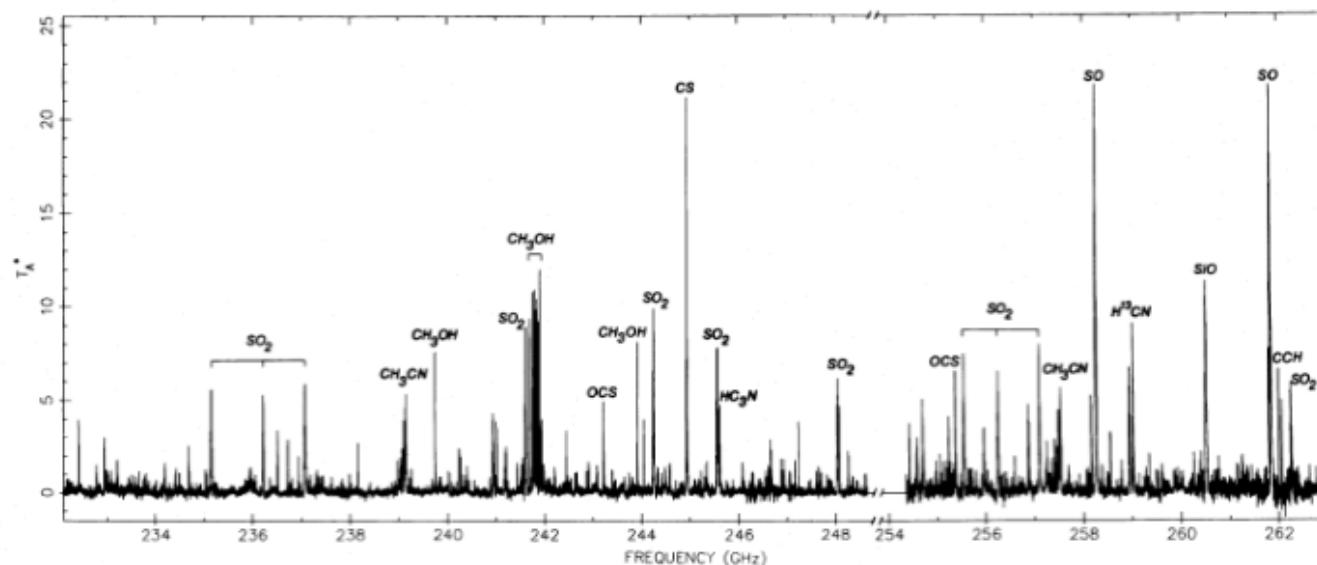
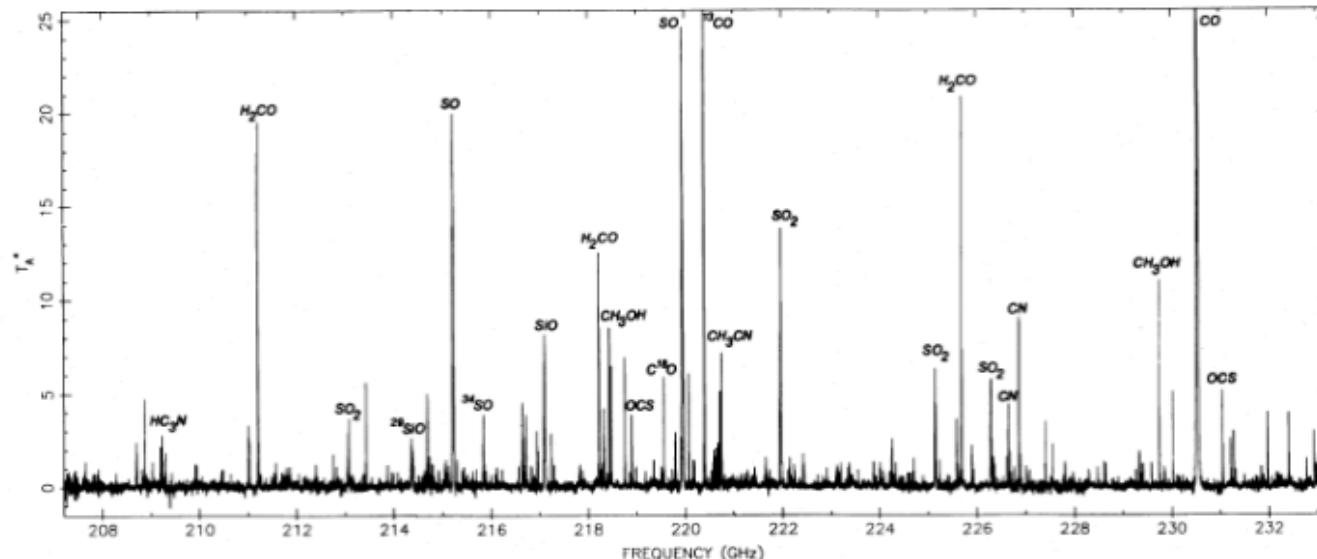
2 atoms	3 atoms	4 atoms	5 atoms	6 atoms	7 atoms	8 atoms	≥ 9 atoms
Hydrogen family							
H ₂	H ₃ ⁺						
Carbon family							
C ₂	C ₃	c-C ₃ H	C ₅	C ₅ H	C ₆ H	C ₇ H	C ₈ H
CH	C ₂ H	l-C ₃ H	C ₄ H	l-C ₄ H ₂	CH ₃ C ₂ H	C ₆ H ₂	CH ₃ C ₄ H
CH ⁺	CH ₂	C ₂ H ₂	l-C ₃ H ₂	C ₂ H ₄		(C ₂ H ₆)	c-C ₆ H ₆
		CH ₂ D ⁺ ?	c-C ₃ H ₂				
			CH ₄				
Oxygen + hydrogen and/or carbon							
CO	C ₂ O	C ₃ O	HCOOH	CH₃OH	HCOCH ₃	HCOOCH₃	(CH ₃) ₂ O
CO ⁺	HCO	HOCO ⁺	H ₂ C ₂ O	HC ₂ CHO	c-C ₂ H ₄ O	CH ₃ COOH	CH ₃ CH ₂ OH
OH	HCO ⁺	H ₂ CO	H ₂ COH ⁺			CH ₂ OHCHO	(CH ₃) ₂ CO
	HOC ⁺	H ₃ O ⁺			CH ₂ CHOH		
	H ₂ O						
	(H ₂ O ⁺)						
	CO ₂						
Nitrogen + hydrogen and/or carbon							
CN	HCN	C ₃ N	HC₃N	CH₃CN	HC ₅ N	CH ₃ C ₃ N	HC ₇ N
NH	HNC	HCCN	CH ₂ CN	CH ₃ NC	CH ₂ CHCN		HC ₉ N
	N ₂ H ⁺	HCN ⁺	HC ₂ NC	HC ₃ NH ⁺	NH ₂ CH ₃		HC ₁₁ N
	NH ₂	H ₂ CN	H ₂ CHN	C ₅ N			CH ₃ CH ₂ CN
		NH ₃	H ₂ NCN				CH ₃ C ₅ N?
			HNC ₃				

2 atoms	3 atoms	4 atoms	5 atoms	6 atoms	7 atoms	8 atoms	≥ 9 atoms
Nitrogen + oxygen + hydrogen and/or carbon							
NO	HNO	HNCO		NH ₂ CHO			NH ₂ CH ₂ COOH?
		N ₂ O					
Molecules with sulphur							
CS	C ₂ S	C ₃ S		CH ₃ SH			
SO	HCS ⁺	HNCS					
SO ⁺	H ₂ S	H ₂ CS					
NS	OCS						
	SO ₂						
Miscellaneous molecules							
SiO	SiCN	SiC ₃	SiC ₄				
SiS	c-SiC ₂		SiH ₄				
HF	MgCN						
SiC	MgNC						
CP	NaCN						
HCl	AlNC						
KCl							
NaCl							
PN							
SiN							
AlF							
AlCl							
SH							
FeO?							

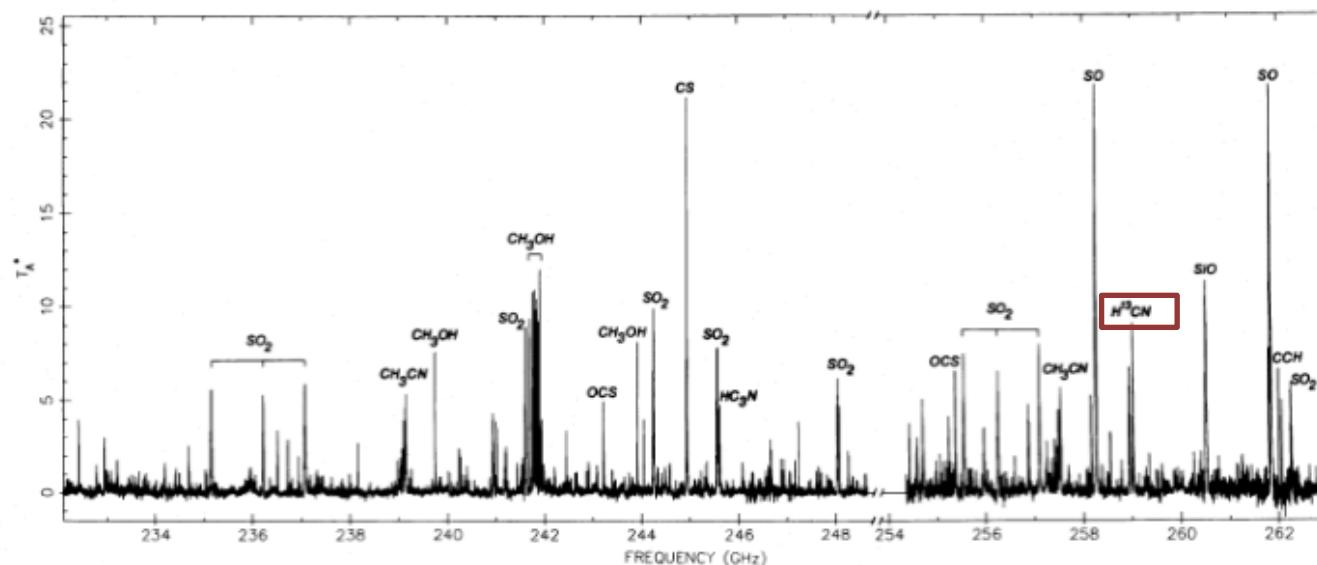
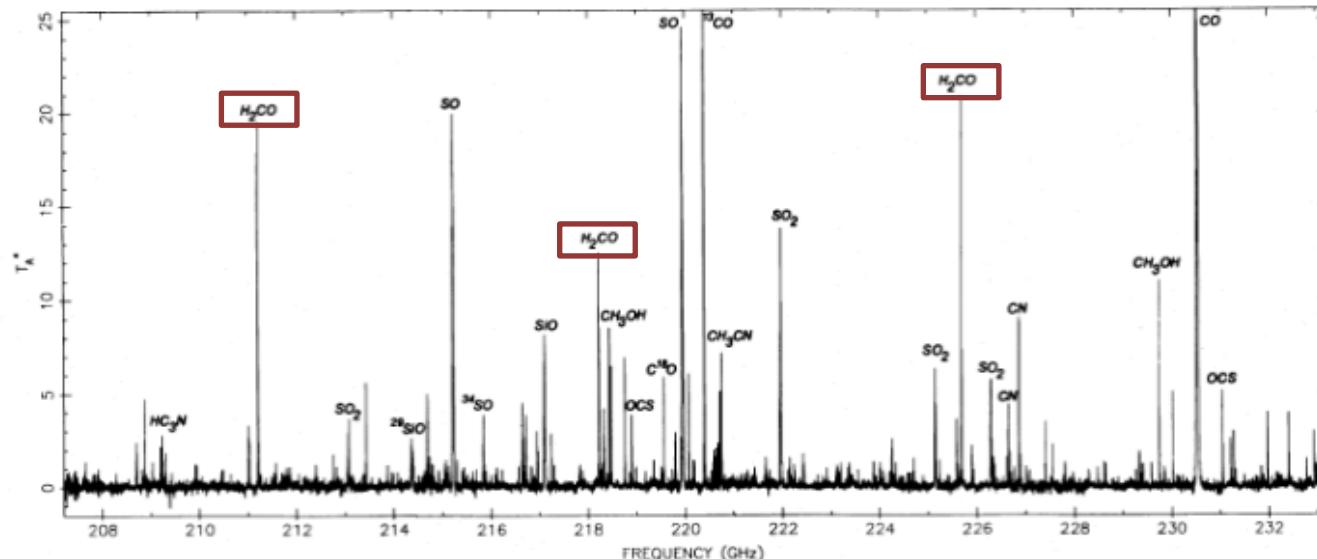
Struktura MC

- Gusto jezgro
- Zbijene strukture
- Centralni oblak molekula
- Fotodisocijativni regioni





Genzel & Stutzki 1989

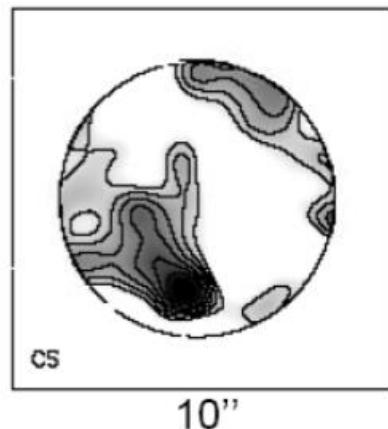
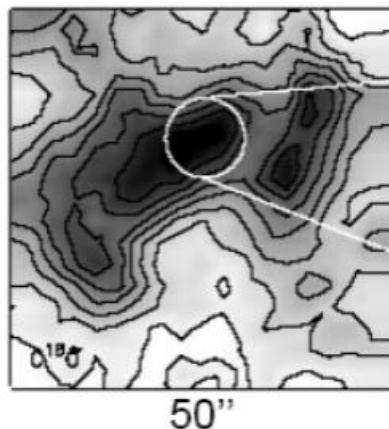
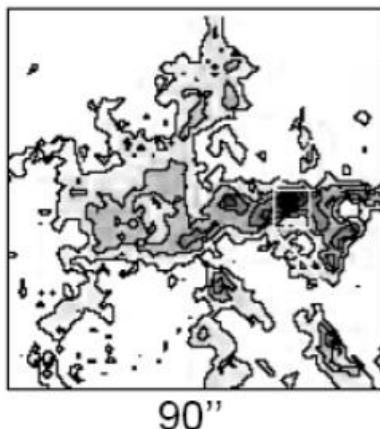


Markeri gustog gasa

Molecule	Transition	Frequency (GHz)	E/k (K)	n_{crit} (cm $^{-3}$) @ 10 K	n_{eff} (cm $^{-3}$) @ 10 K
CS	1-0	49.0	2.4	4.6×10^4	7.0×10^3
	2-1	98.0	7.1	3.0×10^5	1.8×10^4
	3-2	147.0	14	1.3×10^6	7.0×10^4
HCO $^{+}$	1-0	89.2	4.3	1.7×10^5	2.4×10^3
	3-2	267.6	26	4.2×10^6	6.3×10^4
HCN	1-0	88.6	4.3	2.6×10^6	2.9×10^4
	3-2	265.9	26	7.8×10^7	7.0×10^5
H ₂ CO	2 ₁₂ -1 ₁₁	140.8	6.8	1.1×10^6	6.0×10^4
	3 ₁₃ -2 ₁₂	211.2	17	5.6×10^6	3.2×10^5
	4 ₁₄ -3 ₁₃	281.5	30	9.7×10^6	2.2×10^6
NH ₃	(1,1)	23.7	1.1	1.8×10^3	1.2×10^3
	(2,2)	23.7	42	2.1×10^3	3.6×10^4

Struktura MC

- Gusto jezgro
- Zbijene strukture
- Centralni oblak molekula
- Fotodisocijativni regioni



Rotacioni prelazi

- Izraz za energiju: $E_{\text{rot}} = \frac{1}{2I} (\text{angular momentum})^2$
 $(\text{angular momentum})^2 = J(J+1) \left(\frac{\hbar}{2\pi}\right)^2 \quad J = 0, 1, 2, 3, \dots \quad \Delta J = +1 \text{ or } -1$
 - u slučaju linearne molekule $I = \mu r^2$
 - možemo to zapisati na sledeći način:
 $E_{\text{rot}} = J(J+1) \left(\frac{\hbar}{2\pi}\right)^2 \frac{1}{2J} + B = \hbar/(8\pi^2 I) \rightarrow E_{\text{rot}}(J) = \hbar B J(J+1)$
 - ovakav prelazak (u slučaju krutog rotatora) stvara emisiju ili apsorpciju na frekvenciji:

$$\nu = \Delta E / \hbar \Rightarrow B[(J''+1)(J''+2) - J''(J''+1)] = 2B(J''+1)$$

Vibracioni nivoi

- Za odstupanje $r - r_e$ važi: $F = -k(r - r_e)$.
(u klasičnoj mehanici) $E = (1/2)k(r - r_e)^2$.
- U kvantnoj mehanici: $E_{\text{vib}}(v) = \hbar v_{\text{osc}}(v + 1/2)$ $v_{\text{osc}} = 1/(2\pi)\sqrt{(k/\mu)}$.
 $v = 0, 1, 2, 3, \dots$
 - za velike vrednosti v molekuli provode većinu vremena daleko od ravnotežnog položaja
 - privilegovano širenje molekula
 - promena u vrednosti B , odnosno I
 - » *centrifugalna sila*
 $\mu\omega^2r = \mu$ (angular momentum/ I) $^2r = J(J + 1)\hbar^2/(4\pi^2l/(\mu r^3))$
 - » *ravnoteža?*

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- U kvantnoj mehanici: $E_{\text{vib}}(v) = h\nu_{\text{osc}}(v + 1/2)$ $\nu_{\text{osc}} = 1/(2\pi)\sqrt{(k/\mu)}$,
 $v = 0, 1, 2, 3, \dots$
- Usled promena $r \rightarrow$ promena $I = \mu r^2$
 - izraženo za velika v, J (za hladne molekule nije bitno)

$$k(r - r_e) = \frac{J(J+1)\left(\frac{h}{2\pi}\right)^2}{\mu r^3}$$

$$= \frac{J(J+1)\left(\frac{h}{2\pi}\right)^2}{\mu r_e^3 \left(1 + \frac{r - r_e}{r_e}\right)^3}$$

$$\begin{aligned} \text{total energy} &= \frac{J(J+1)\left(\frac{h}{2\pi}\right)^2}{2\mu r^2} - \frac{1}{2}k(r - r_e)^2 \\ &= \frac{J(J+1)\left(\frac{h}{2\pi}\right)^2}{2\mu r_e^2 \left(1 + \frac{r - r_e}{r_e}\right)^2} - \frac{1}{2}k(r - r_e)^2 \\ &\simeq hBJ(J+1) - \frac{J^2(J+1)^2h^4}{32\pi^4\mu^2r_e^6k} \end{aligned}$$

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$$\boxed{(r - r_e)/r_e \ll 1}$$

$$k(r - r_e) = \frac{J(J+1)\left(\frac{h}{2\pi}\right)^2}{\mu r^3}$$

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$$E = hBJ(J+1) - hDJ^2(J+1)^2 + \dots$$

$$D = 16\pi^2 B^3 \mu/k = 4B^3/v_{\text{osc}}^2$$

$$v = 2B(J''+1) - 4D(J''+1)^3$$

$$B_v = B_e - \alpha(v + 1/2) + \dots \quad D_v = D_e - \beta(v + 1/2)$$

Verovatnoća?

- verovatnoća prelaza je data sa: $A_{ji} = \frac{64\pi^4}{3hc^34\pi\epsilon_0} v^3 |d|^2$
 - dipolni momenat: $d = e \int \Phi_j r \Phi_i dV$
$$d^2 = f(J) (e \int \phi_j r \phi_i dV)^2 = \frac{(J+1)}{(2J+3)} d_p^2$$
 - verovatnoća:
$$A_{ji} = 1.046 * 10^{21} v^3 d_p^2 \frac{(J+1)}{(2J+3)}$$
 - particiona funkcija:

$$U = \sum_j (2J+1) \exp \left[-\frac{hBJ(J+1)}{kT_{ex}} \right] \simeq \int_0^\infty (2J+1) \exp \left[-\frac{hBJ(J+1)}{kT_{ex}} \right] dJ$$
$$\simeq \int_0^\infty \exp \left[-\frac{hBJ'}{kT_{ex}} \right] dJ', \text{ with } J' = J^2 + J$$
$$\simeq \frac{kT_{ex}}{hB}$$

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 - verovatnoća:
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 - particiona funkcija:
$$U = (kT_{ex})/(hB) + 1/3$$

$$\begin{aligned} U &= \sum_j (2J+1) \exp \left[-\frac{hBJ(J+1)}{kT_{ex}} \right] \simeq \int_0^\infty (2J+1) \exp \left[-\frac{hBJ(J+1)}{kT_{ex}} \right] dJ \\ &\simeq \int_0^\infty \exp \left[-\frac{hBJ'}{kT_{ex}} \right] dJ', \quad \text{with } J' = J^2 + J \\ &\simeq \frac{kT_{ex}}{hB} \end{aligned}$$

Verovatnoća?

- verovatnoća prelaza je data sa:

CO u osnovnom stanju
nema nuklearni spin
a ugaoni momenat = 0



$$U = \sum_j ($$

$$\simeq \int_0^\infty \exp\left[-\frac{hBJ'}{kT_{\text{ex}}}\right] dJ', \quad \text{with } J' = J^2 + J$$

$$\simeq \frac{kT_{\text{ex}}}{hB}$$

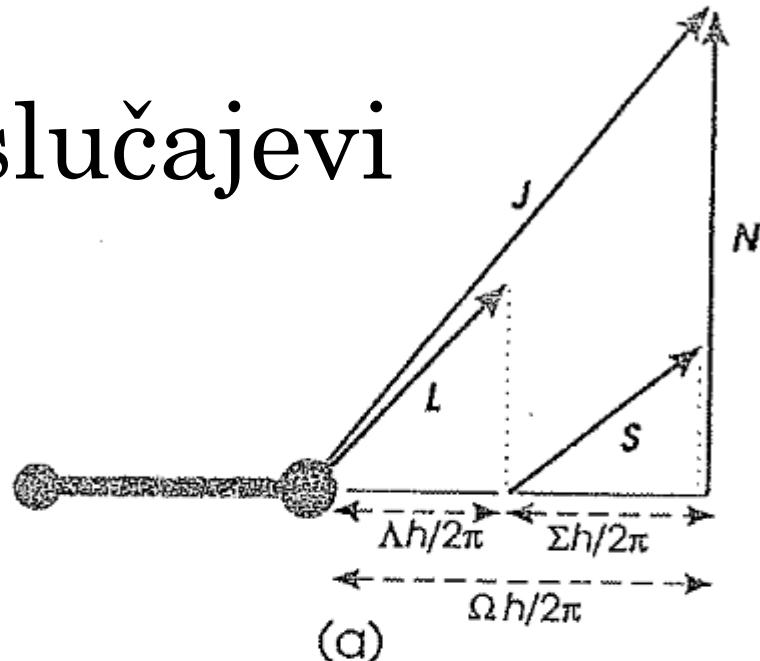
$$\begin{aligned} 1) & d_p^2 \\ 3) & \end{aligned}$$

Šta utiče na energiju i intenzitet rotacionih i vibracionih linija?

dužina molekula, svedena masa
dipolni momenat, ose simetrije
superpozicija kvantnih stanja,
prostorna orijentacija molekula...

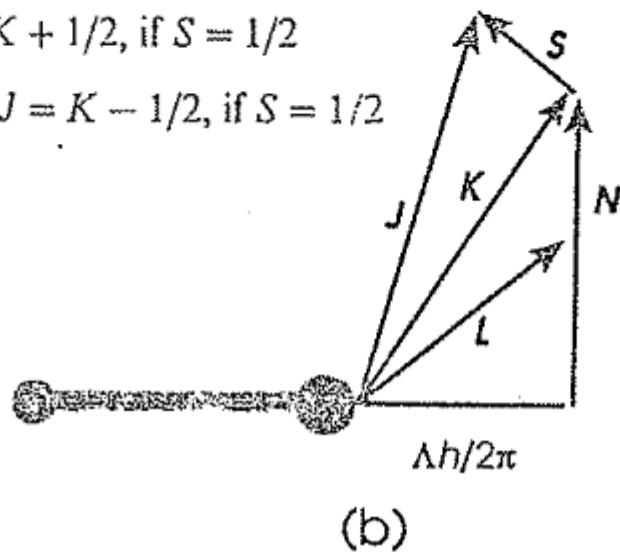
Hundovi slučajevi

- Linearan molekul
- Slaganje spinova



$$\text{energy} = hBK(K+1) + \text{constant}*K \quad \text{for } J = K + 1/2, \text{ if } S = 1/2$$

$$\text{energy} = hBK(K+1) - \text{constant}*(K+1) \quad \text{for } J = K - 1/2, \text{ if } S = 1/2$$



Ne-linearni molekuli

- Uopšten slučaj

$$E = \frac{1}{2}I_x\omega_x^2 + \frac{1}{2}I_y\omega_y^2 + \frac{1}{2}I_z\omega_z^2 = \frac{(\text{angular Momentum about } x)^2}{2I_x} + \frac{(\text{angular Momentum about } y)^2}{2I_y} + \frac{(\text{angular Momentum about } z)^2}{2I_z}$$

- U sistemu gde je jedna od osa odabrana za osu molekula:

$$E = \hbar BJ(J+1) + \hbar(A-B)\Omega^2$$

$$A = \hbar/(8\pi^2 I_z)$$

$$(\text{angular momentum about } x)^2 + (\text{angular momentum about } y)^2$$

$$= N^2 = J^2 - \Omega^2 = (J(J+1) - \Omega^2) [h/(2\pi)]^2$$

- Kvadripolni član takođe utiče na linije
 - za izražen magnetni momenat

Šta utiče na energiju i intenzitet rotacionih i vibracionih linija?

dužina molekula, svedena masa
dipolni momenat, ose simetrije
superpozicija kvantnih stanja,
prostorna orijentacija molekula...

Oblik molekula

- Momenti inercija duž glavnih osa su obeleženi u skladu sa konvencijom:

$$I_c \geq I_b \geq I_a$$

Definicije

- Simetrični
 - elektronski prelasci dozvoljeni
 - rotacioni i vibracioni nisu

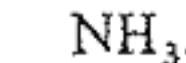
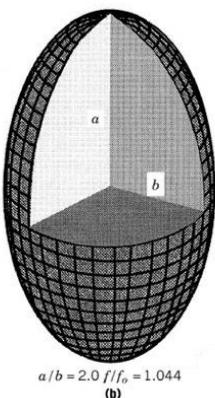
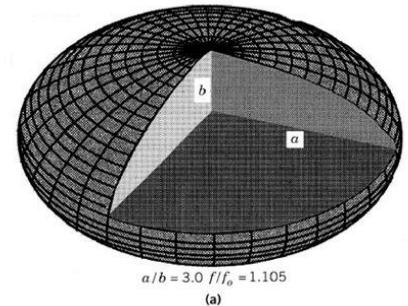
- Sferno simetrični

- Kružno simetrični (Symmetric Top)

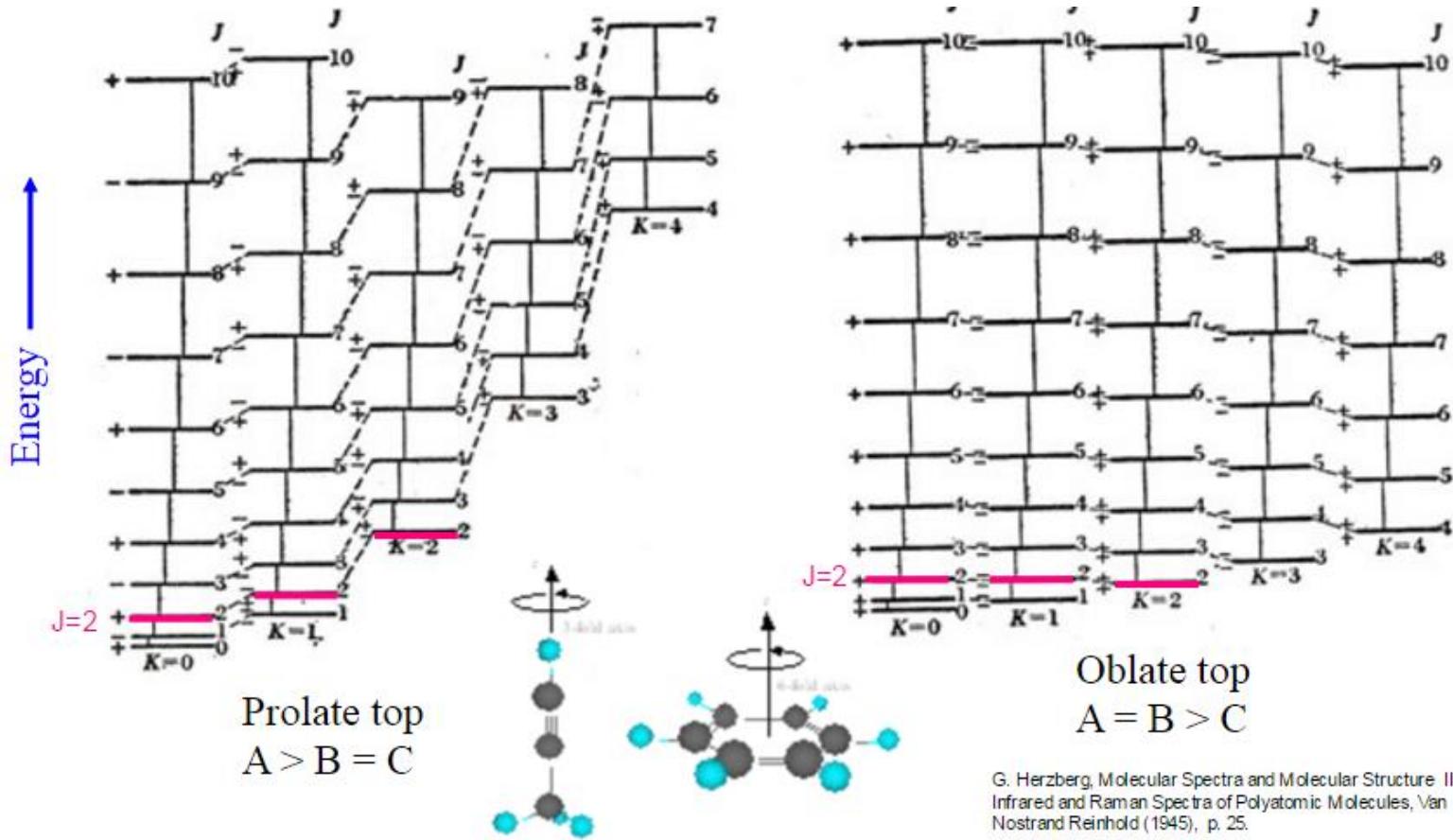
– Oblate (a) $I_c > I_b = I_a$

– Prolate (b) $I_c = I_b > I_a$

- Asimetrični



Spektar Prolate/Oblate molekula



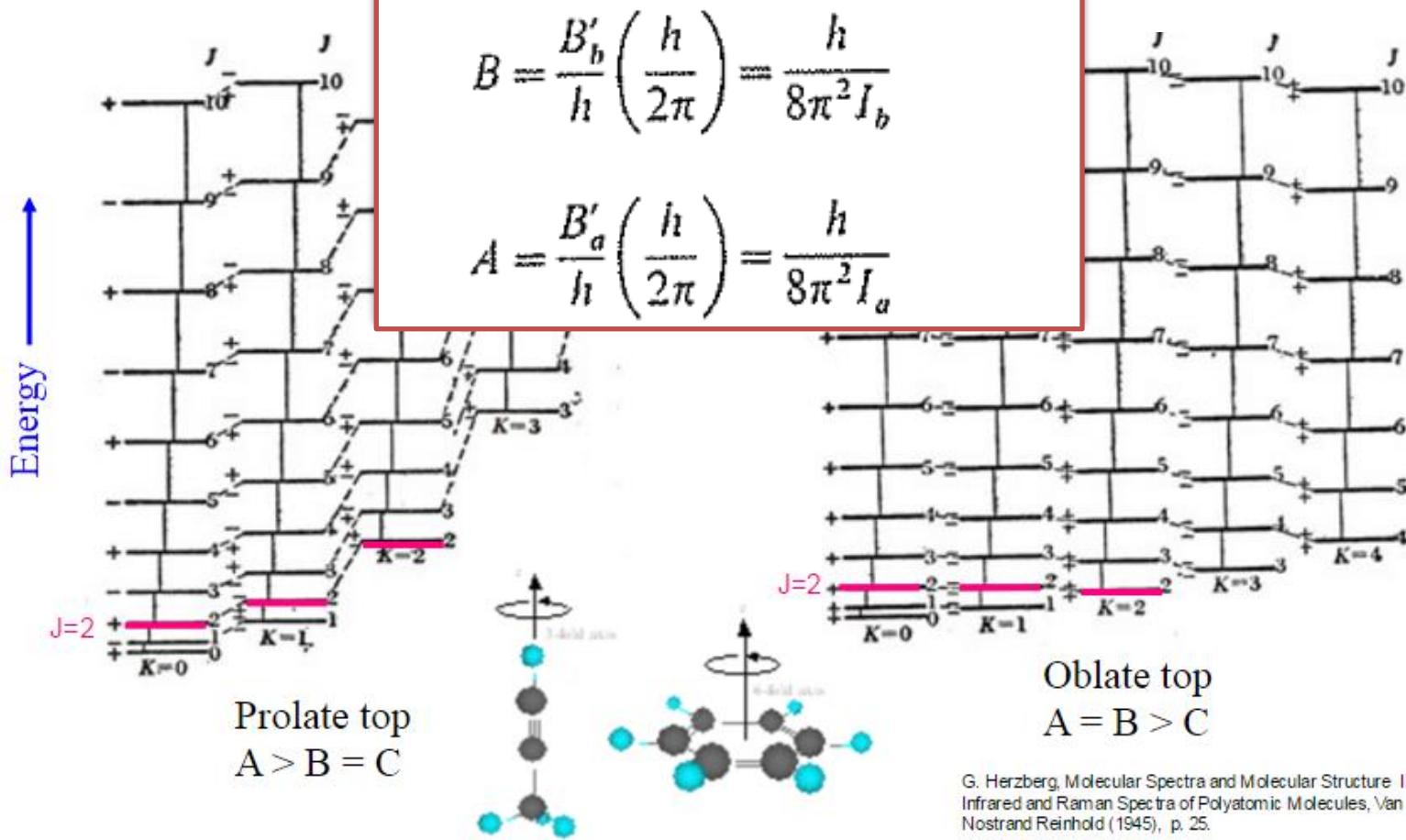
G. Herzberg, Molecular Spectra and Molecular Structure II. Infrared and Raman Spectra of Polyatomic Molecules, Van Nostrand Reinhold (1945), p. 25.

Oblate: benzene, C_6H_6 , ammonia, NH_3

Prolate: chloromethane, CH_3Cl , propyne, $CH_3C\equiv CH$

Spektren von Molekülen

$$E(J, K) = hBJ(J + 1) + h(A - B)K^2$$



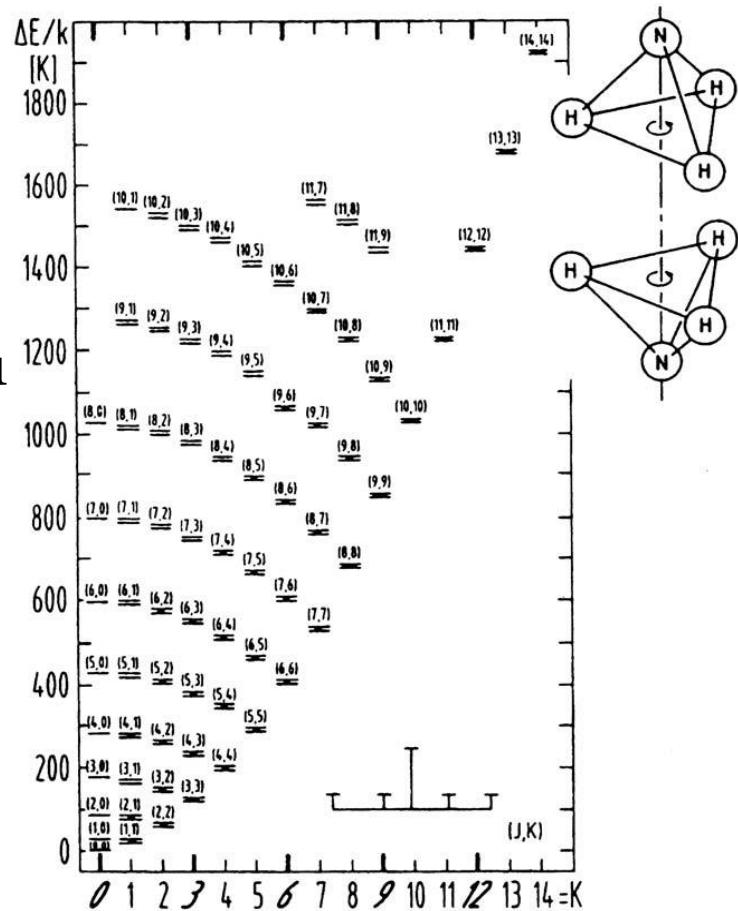
G. Herzberg, Molecular Spectra and Molecular Structure II. Infrared and Raman Spectra of Polyatomic Molecules, Van Nostrand Reinhold (1945), p. 25.

Oblate: benzene, C_6H_6 , ammonia, NH_3

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Orto i Para forme

- Orto – paralelni nuk. spin
- Para – anti-paralelan spin
- Prelazak iz jedne u drugu formu je izričito zabranjen prelaz
 - detekcija H_2
 - temperatura formiranja molekula



Jačina linija molekula – primer na CO

- Kako odrediti
 - Temperaturu?
 - Linijsku gustinu?
 - CO ima više prelaza i može biti sačinjen od više izotopa
 - niz linija
 - raspodela po energetskim nivoima
 - potvrđuju/opovrgavaju LTE
 - daju uvid u T sredine
 - optički guste/tanke
 - niz frekvencija u radio, IC itd.
 - CMB ili drugi pozadinski izvor

Temperatura sjaja

$$\begin{aligned}\int T_B dv &= \frac{c^2 h}{8\pi k} \frac{A_{ji} g_j}{g_i} \frac{1}{v} \frac{g_i}{U} \exp\left[-\frac{E}{kT_{ex}}\right] (N_{TOT} \rho R) \\ &= \frac{c^2 h}{8\pi k} \frac{A_{ji}(2J+3)}{(2J+1)} \frac{1}{2B(J+1)} \frac{hB}{kT_{ex}} (2J+1) \exp\left[-\frac{hBJ(J+1)}{kT_{ex}}\right] (N_{TOT} \rho R) \\ &= \frac{c^2 h^2}{16\pi k^2} \frac{1}{T_{ex}} A_{ji} \frac{(2J+3)}{(J+1)} \exp\left[-\frac{hBJ(J+1)}{kT_{ex}}\right] (N_{TOT} \rho R)\end{aligned}$$

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$$\begin{aligned}\int T_B dv &= \frac{4\pi^3 h}{3ck^2(4\pi\varepsilon_0)} \frac{v^3}{T_{ex}} d_p^2 \exp\left[-\frac{hBJ(J+1)}{kT_{ex}}\right] (N_{TOT} \rho R) \\ &= \frac{32\pi^3 h}{3ck^2(4\pi\varepsilon_0)} \frac{B^3(J+1)^3}{T_{ex}} d_p^2 \exp\left[-\frac{hBJ(J+1)}{kT_{ex}}\right] (N_{TOT} \rho R)\end{aligned}$$

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Koliko je optički gusta sredina?

$$\begin{aligned}\int T_B dv &= \frac{4\pi^3 h}{3ck^2(4\pi\varepsilon_0)} \frac{v^3}{T_{ex}} d_p^2 \exp\left[-\frac{hBJ(J+1)}{kT_{ex}}\right] (N_{TOT} \rho R) \\ &= \frac{32\pi^3 h}{3ck^2(4\pi\varepsilon_0)} \frac{B^3(J+1)^3}{T_{ex}} d_p^2 \exp\left[-\frac{hBJ(J+1)}{kT_{ex}}\right] (N_{TOT} \rho R)\end{aligned}$$

Linije različitih izotopa

- Odnos između nekoliko izotopa C

$$N(^{12}\text{CO}) \simeq 76 N(^{13}\text{CO})$$

$$N(^{12}\text{CO}) = 560 N(\text{C}^{18}\text{O})$$

- Za dugo-talasnu aproksimaciju ove linije imaju istu:
 - temperaturu pozadinskog zračenja $T_B = T_{\text{ex}}(1 - e^{-\tau})$
 - temperaturu eksicitacije T_{ex}

- Stoga važi: $\frac{T_B(^{12}\text{C}^{16}\text{O})}{T_B(^{13}\text{C}^{16}\text{O})} = \frac{(1 - e^{-\tau(^{12}\text{C}^{16}\text{O})})}{(1 - e^{-\tau(^{13}\text{C}^{16}\text{O})})}$

$$N(^{12}\text{C})/N(^{13}\text{C}) = R,$$

$$\tau(^{12}\text{C}^{16}\text{O}) = R \tau(^{13}\text{C}^{16}\text{O})$$

optički tanak sl.

$$\tau = T_B/T_{\text{ex}}$$

$$\int \tau_v dV = \frac{8\pi^3 B d_p^2}{3k4\pi\varepsilon_0} \frac{N_{\text{TOT}} \rho R(J+1) e^{-hBJ(J+1)/kT_{\text{ex}}}}{T_{\text{ex}} + \frac{hB}{3k}} (1 - e^{-hv/kT_{\text{ex}}})$$

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optički tanak sl.
 $\tau = T_B/T_{\text{ex}}$

$$N_{\text{TOT}} \rho R(^{12}\text{C}^{16}\text{O}) = 2.39 \times 10^{15} \frac{e^{hB_{12}J(J+1)/kT_{\text{ex}}}}{J+1} \frac{T_{\text{ex}} + 0.93}{1 - e^{-hv_{12}/hT_{\text{ex}}}} \int \tau dV \text{ m}^{-2}$$

$$N_{\text{TOT}} \rho R(^{13}\text{C}^{16}\text{O}) = 2.49 \times 10^{15} \frac{e^{hB_{13}J(J+1)/kT_{\text{ex}}}}{J+1} \frac{T_{\text{ex}} + 0.89}{1 - e^{-hv_{13}/hT_{\text{ex}}}} \int \tau dV \text{ m}^{-2}$$

Prelazi na različit nivo

- Unutar istog molekula od istih izotopa posmatramo prelaska na (niži) nivo l, m

$$\frac{\tau_V(m \text{ to } m+1)}{\tau_V(l \text{ to } l+1)} = \frac{(J_m + 1) e^{-E_m/kT_{\text{ex}}}}{(J_l + 1) e^{-E_l/kT_{\text{ex}}}} \frac{1 - e^{-hv_m/kT_{\text{ex}}}}{1 - e^{-hv_l/kT_{\text{ex}}}}$$

- Ukoliko su nivoi uzastopni ($m=l+1$)

$$\frac{\tau_V(l+1 \text{ to } l+2)}{\tau_V(l \text{ to } l+1)} = \frac{J_l + 2}{J_l + 1} \frac{1 - e^{-hv_{m+1}/kT_{\text{ex}}}}{e^{hv_l/kT_{\text{ex}}} - 1}$$

Detaljnije – sa pozadinskim zračenjem

$$T'_B = c^2/(2k\nu^2)I_\nu$$

- Prisutnost pozadinskog zračenja ($T=2.7\text{K}$)

$$T_B = T_{\text{ex}}(1 - e^{-\tau}) + T_{\text{BCK}}e^{-\tau} \quad \text{or} \quad T_{\text{OBS}} = T_B - T_{\text{BCK}} = (T_{\text{ex}} - T_{\text{BCK}})(1 - e^{-\tau})$$

- Ukoliko predstavimo preko Planka

$$I_\nu = \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/kT_{\text{BCK}}} - 1} e^{-\tau_\nu} + \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/kT_{\text{ex}}} - 1} (1 - e^{-\tau_\nu})$$

$$T'_{\text{OBS}} = T'_B - T'_{\text{BCK}} = \frac{h\nu}{k} \left[\frac{1}{e^{h\nu/kT_{\text{ex}}} - 1} - \frac{1}{e^{h\nu/kT_{\text{BCK}}} - 1} \right] (1 - e^{-\tau_\nu})$$

$$T'_{\text{OBS}}(1-0) = 5.532 \left(\frac{1}{[e^{5.532/T_{\text{ex}}} - 1]} - 0.1479 \right) (1 - e^{-\tau_\nu}) \quad \text{for } {}^{12}\text{C}{}^{16}\text{O}$$

$$= 5.289 \left(\frac{1}{[e^{5.289/T_{\text{ex}}} - 1]} - 0.1642 \right) (1 - e^{-\tau_\nu}) \quad \text{for } {}^{13}\text{C}{}^{16}\text{O}$$

Rezultat razmatranja

- Ako uzmemo da je temperatura ekscitacije nekoliko puta veća od pozadinske dobijamo:

$$N_{\text{TOT}} \rho R = 4.3 \times 10^{14} \frac{(T_{\text{ex}} + 0.93)}{e^{-5.53/T_{\text{ex}}}} \int \frac{T'_{\text{OBS}} \tau_\nu}{(1 - e^{-\tau_\nu})} dV$$

Molekul sa dva nivoa

- Nalazi se u polju zračenja J_v sa T_B

$$S_v = \frac{2hv^3}{c^2} \frac{1}{e^{hv/kT_{ex}} - 1} = \frac{\int \phi_v J(T_R)_v + \varepsilon B_v(T_K)}{1 + \varepsilon}$$

$$\begin{aligned}\varepsilon &= \frac{C_{ji}}{A_{ji}} (1 - e^{-hv/kT_K}) \\ &\quad \xrightarrow{\text{blue arrow}} \varepsilon = \frac{C_{ji}}{A_{ji}} \frac{hv}{kT_K}\end{aligned}$$

- Uzimamo nekoliko aprox.

$$\int \phi_v J_v dv = J_v \int \phi_v dv = J_v(T_R)$$

$$S_v = (2kv^2)/c^2 T_{ex}$$

$$J_v = (2kv^2)/c^2 T_R$$

$$B_v = (2kv^2)/c^2 T_K$$

- Dobijamo: $T_{ex} = \frac{T_R + \varepsilon T_K}{1 + \varepsilon}$

- Odnosno...

$$T_{ex} = T_K \left(\frac{T_0 + T_B}{T_0 + T_K} \right) \quad T_0 = \frac{hv}{k} \frac{C_{ul}}{A_{ul}}$$

Razmatranje molekula sa dva nivoa

- Kakav je oblik jednačine koju smo dobili?

$$T_{ex} = T_K \left(\frac{T_0 + T_B}{T_0 + T_K} \right) \quad T_0 = \frac{h\nu}{k} \frac{C_{ul}}{A_{ul}}$$

- Ukoliko su sudari dominantni:

- T_0 je veliko i $T_{ex} \sim T_K$ (ravnoteža sudarima/LT)

- Ukoliko sudari nisu dovoljno česti:

- T_0 je malo i $T_{ex} \sim T_B$ (ravnoteža zračenja)

- Bez dugotalasne aproksimacije

$$\frac{1}{e^{h\nu/kT_{ex}} - 1} = \frac{1}{1 + \varepsilon} \frac{1}{[e^{h\nu/kT_K} - 1]} + \frac{\varepsilon}{1 + \varepsilon} \frac{1}{[e^{h\nu/kT_B} - 1]}$$

$$\frac{T_0}{T_{ex}} = \frac{T_0}{T_K} + \ln_e \left(\frac{\frac{e^{T_0/T_K} + \frac{C_{ji}}{A_{ji}} [e^{T_0/T_K} - 1]}{e^{T_0/T_K} + \frac{C_{ji}}{A_{ji}} [e^{T_0/T_K} - 1]}}{\frac{e^{T_0/T_B} + \frac{C_{ji}}{A_{ji}} [e^{T_0/T_B} - 1]}{e^{T_0/T_B} + \frac{C_{ji}}{A_{ji}} [e^{T_0/T_B} - 1]}} \right)$$

Pojam kritične gustine

Procesi u MC

- **Hladjenje gasa**
 - Rotacioni prelazi i emisione linije CO i drugih molekula
 - Sudari molekla sa hladnim česticama prešine
 - Emisija sudarno pobudenih linija [C ii], [Si ii] [C i]
 - Rotaciono-vibracione linje molekulskog vodonika
- **Zagrevanje gasa** (OB zvezde, difuzno polje zračenja)
 - Fotoelektricni mehanizam zagrevanja
 - UV-fluorescentno zagrevanje (Ly-Weiber fotoni)

Određivanje mase MC

- Džinsov zakon daje:

$$\mathcal{M}_J = \frac{4\pi}{3} R_J^3 \rho_c \approx 18 \mathcal{M}_\odot T^{1.5} n^{-0.5}$$

- Na osnovu teoreme virijala:

$$\mathcal{M}_{\text{vir}} \approx \Delta v^2 \left(\frac{R_c}{G} \right) \quad \Delta v = \sqrt{\frac{GM_c}{R_c}}$$

- Na osnovu posmatranja CO:

$$L_{\text{CO}} = D^2 \int I_{\text{CO}} d\Omega \Leftrightarrow L_{\text{CO}} = T_A^*(\text{CO}) \Delta v \pi R_c^3$$

- **Konačna relacija:**

$$L_{\text{CO}} = T_A^*(\text{CO}) \left(\frac{3\pi G}{4\rho_c} \right)^{1/2} \mathcal{M}_c$$

Odnos vodonika i CO?

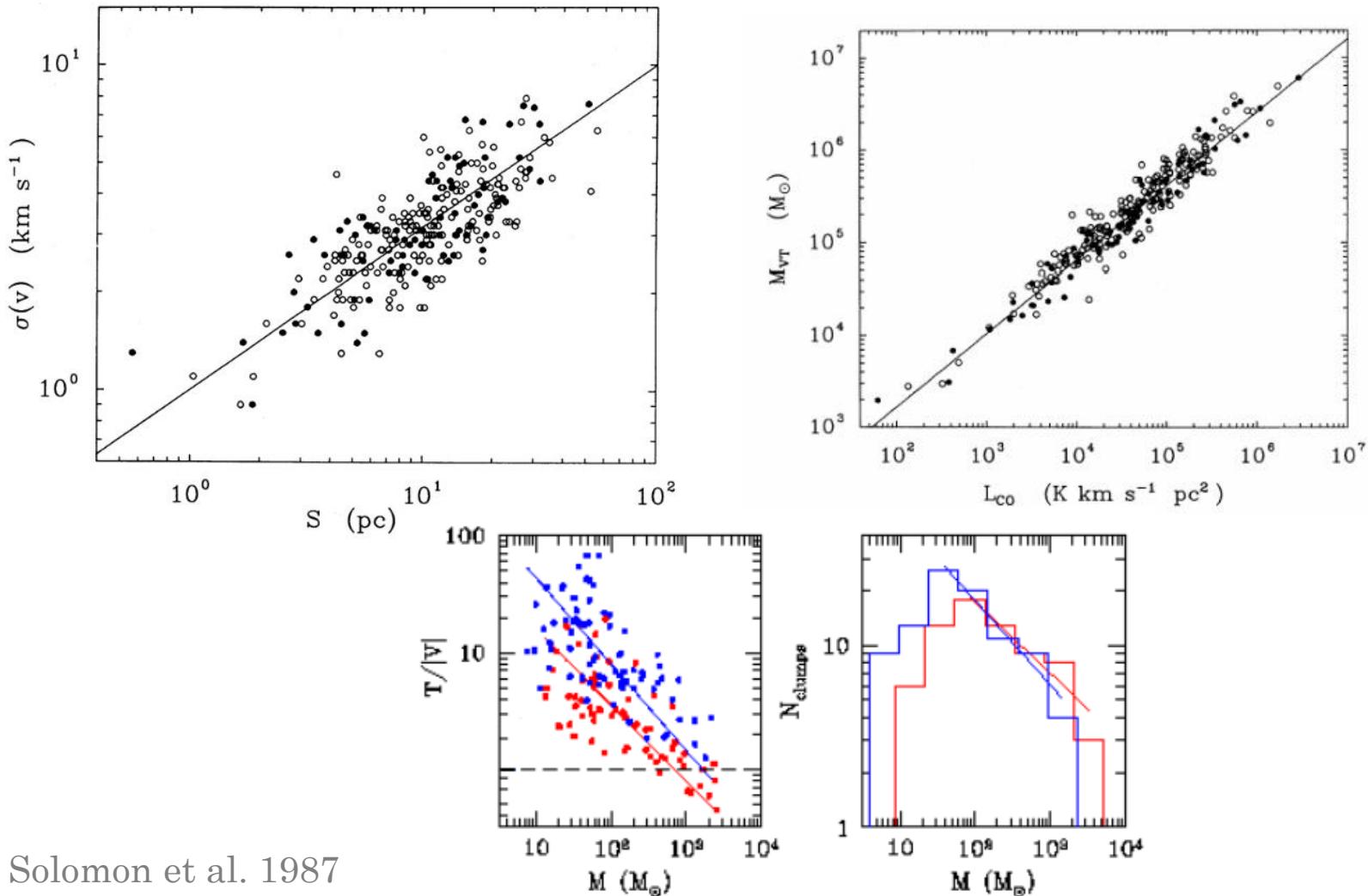
- Odnos se smatra konstantnim – određen je **X faktorom** u skladu sa formulom:

$$\frac{N(\text{H}_2)}{I_{\text{CO}}} = X \cdot 10^{20} \text{ cm}^{-2} [\text{K km s}^{-1}]$$

- Vrednosti dobijene za X variraju

Early work	2-5
γ -rays Hunter et al. 1997	1.56 ± 0.05
HI/IRAS/CO Dame et al. 2001	1.8 ± 0.3
A_V Lada et al. 2003	~ 4

Emprijski zakoni



Ključna pitanja

Ostaje nekoliko važnih tema koje nisu razjašnjene

- *Da li se GMC formiraju u spiralnim granama galaksija?*
- *Koji od svih procesa (sudari manjih molekulskih struktura, nestabilnost odnosa gravitacija - termalni gas, galaktički udarni talasi/najčešće supernova ostaci, magnetna polja) je najdominantniji u stvaranju molekulskih oblaka?*
- *Da li su sve površinske gustine u GMC iste ili bar približno iste?*
- *Možemo li strukturu GMC-a uopštiti kao diskretnu?*
- *Da li postoji veza izmedu fenomena dugog „životnog veka“ spiralnih galaksija sa eventualnim dugim vremenom života molekulskih oblaka koji su zastupljeni u njima?*